#### "Finding" a Pulse Shape

Jason Taylor May 16, 2005

"The laser cavity finds the pulse that minimizes loss--it's like magic."

## My Research

I investigate the usefulness of laser cavities and laser dynamics for information processing.

Kerr Lens Mode-Locked laser cavities

## **Project Goals**

- Discover what a KLM laser minimizes
  - Cavity loss?
  - Population Inversion?
- Use results to predict good bit representation and/or logical operators
  - Space
  - Time
  - Phase
  - Power

#### Kerr-Lens Mode Locked Oscillator



## The Equation



This equation describes the evolution of a short pulse over one round trip in a KLM cavity.



## Mode Locking



#### Artificial Fast Saturable Absorbers



[Hau00]



master equation:

$$\frac{1}{T_R}\frac{\delta}{\delta T}a = (g-l)a + (\frac{g}{\Omega_g^2} + \frac{1}{\Omega_f^2})\frac{\delta^2}{\delta t^2}a + \gamma |a|^2a$$

 $a_0(t) = A_0 \operatorname{sech}(t/\tau)$  solution unbounded

Siegman, Haus [Hau00]

#### Kerr-Lens Mode Locked Laser



Ti:Sapphire Absorption/Emission Spectra

Ti:Sapphire KLM oscillators are available commercially—where else would this graphic come from?

700

Wavelength (Nanometers)

800

900

1000

400

500

600

#### Soliton Effects in Ultrashort Pulses

 $\Delta n = n_2 I(t)$ 

Self Phase Modulation (SPM)

$$\Delta a = -j\delta |a|^2 a$$

Group Velocity Dispersion (GVD)

$$\Delta a = jD \frac{d^2}{dt^2} a$$



$$\frac{1}{T_R}\frac{\delta}{\delta T}a = (g-l)a + \left(\frac{1}{\Omega_f^2} + jD\right)\frac{\delta^2}{\delta t^2}a + (\gamma - j\delta)|a|^2a$$

$$\frac{1}{GVD}\frac{\delta^2}{\delta T}a = (g-l)a + \left(\frac{1}{\Omega_f^2} + jD\right)\frac{\delta^2}{\delta t^2}a + (\gamma - j\delta)|a|^2a$$

$$\delta_0 = \text{no SPM}$$
  
 $D_n = \text{GVD}$ 

#### Haus' Master Equation

master equation with GVD and SPM:

$$\frac{1}{T_R}\frac{\delta}{\delta T}a = (g-l)a + \left(\frac{1}{\Omega_f^2} + jD\right)\frac{\delta^2}{\delta t^2}a + (\gamma - j\delta)|a|^2a$$

$$\frac{1}{GVD}$$

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Gain depletion

$$g = g(T) = \frac{g_0}{1 + \frac{E_P(T)}{E_{sat}}}$$

where

$$E_P(T) = \int_{-\infty}^{+\infty} |A(T,t)|^2 dt$$
$$E_P(T) = \Delta t \sum_{n=0}^{N} |A(T,n)|^2$$

#### Numerical Simulation

## Here we look at the equation properties via numerical simulation.

Separate into Linear and Non-Linear Operators

$$\frac{1}{T_R}\frac{\delta}{\delta T}a = (g-l)a + \left(\frac{1}{\Omega_f^2} + jD\right)\frac{\delta^2}{\delta t^2}a + (\gamma - j\delta)|a|^2a$$

$$\frac{\partial A}{\partial T} = (\hat{D} + \hat{N})A$$

where

$$\hat{D} = g - l + (D_{gf} + jD) \frac{\partial^2}{\partial t^2}$$
$$\hat{N} = (\gamma - j\delta) |A|^2$$

 $A(T + \kappa, t) = \exp[\kappa(\hat{D} + \hat{N})]A(T, t)$ 

#### Simulation



#### Gain, SAM, no GVD, no SPM





## What is minimized?

- Cavity loss?
- Gain medium population inversion?
- Pulse width?

#### Haus' Master Equation

master equation with GVD and SPM:

$$\frac{1}{T_R}\frac{\delta}{\delta T}a = (g-l)a + \left(\frac{1}{\Omega_f^2} + jD\right)\frac{\delta^2}{\delta t^2}a + (\gamma - j\delta)|a|^2a$$

$$\frac{1}{GVD}$$

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#### Complex Ginzburg-Landau Equation

master equation with GVD and SPM:

$$\frac{1}{T_R}\frac{\delta}{\delta T}a = (g-l)a + \left(\frac{1}{\Omega_f^2} + jD\right)\frac{\delta^2}{\delta t^2}a + (\gamma - j\delta)|a|^2a$$

$$\frac{1}{GVD}$$

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Gain depletion

$$g = g(T) = \frac{g_0}{1 + \frac{E_P(T)}{E_{sat}}}$$

where

$$E_P(T) = \int_{-\infty}^{+\infty} |A(T,t)|^2 dt$$
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#### Complex Ginzburg-Landau Equation

master equation:

$$\frac{1}{T_R}\frac{\delta}{\delta T}a = (g-l)a + \left(\frac{1}{\Omega_f^2} + jD\right)\frac{\delta^2}{\delta t^2}a + (\gamma - j\delta)|a|^2a$$

soliton like pulse

CW solution

$$a(t) = A_0 \operatorname{sech}^{(1+j\beta)}\left(\frac{t}{\tau}\right)$$

$$a(t) = A_0 \exp(-j\omega t)$$

#### general CGLE

$$\frac{\partial}{\partial T}A = A + (1 + jc_1)\frac{\partial^2}{\partial t^2}A - (1 + jc_2)|A|^2A$$

# Complex Ginzburg-Landau Equation general CGLE

$$\frac{\partial}{\partial T}A = A + (1 + jc_1)\frac{\partial^2}{\partial t^2}A - (1 + jc_2)|A|^2A$$



#### Lyapunov Function

soliton like pulse

CW solution

$$a(t) = A_0 \operatorname{sech}^{(1+j\beta)}\left(\frac{t}{\tau}\right)$$

$$a(t) = A_0 \exp(-j\omega t)$$

## A good approximate Lyapunov function is known for a CW stationary solution.

No Lyapunov function is known for soliton solutions—too close to chaos. IMHO

## Simplify Equation

$$\frac{1}{T_R}\frac{\delta}{\delta T}a = (g-l)a + \left(\frac{1}{\Omega_f^2} + jD\right)\frac{\delta^2}{\delta t^2}a + (\gamma - j\delta)|a|^2a$$

$$\frac{\partial}{\partial T}A(T,t) = jD\frac{\partial^2}{\partial t^2}A - j\delta|A|^2A$$

ignore gain depletion, BW filtering and SAM

#### Non-linear Schrodinger Equation

$$\frac{\partial}{\partial T}A(T,t) = jD\frac{\partial^2}{\partial t^2}A - j\delta|A|^2A$$

Lyapunov Function:

$$V = \int_{-\infty}^{+\infty} dt \left[ -D|A|^2 + \frac{\delta}{4}|A|^4 a + \left| \frac{\partial}{\partial t} A\right|^2 \right]$$

Looks a lot like minimizing the action.

#### Numerical Confirmation

#### Try Lyapunov function in simulator.

#### NLSE Lyapunov function on CGLE with gain saturation





#### Zoom in on Hump



#### Zoom in on 2<sup>nd</sup> Hump



#### Zoom in on 3<sup>rd</sup> Hump



#### Zoom in on 4<sup>th</sup> Hump



## Conclusions

- We found a fractal.
- The NLSE approximate Lyapunov function isn't valid far away from the soliton solution.
- Is a numerical stability analysis of the CGLE sufficient?